

A New Characterization of Euler Elastica

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Elastic Behavior of Roads and Beams

- Galileo Galilei (around 1638) asked the question about the force required to break a beam set into a wall.
- James Bernoulli (1687-1692) raised the question of the shape of a beam.
- James Bernoulli (1694) published the first solution of the rectangular elastica.
- Daniel Bernoulli (1742) proposed variational techniques for the problem of the elastica.
- Leonhard Euler (1743) solved the general problem and classified the elastica shapes.

Euler Elastica

A Brief Historical Overview

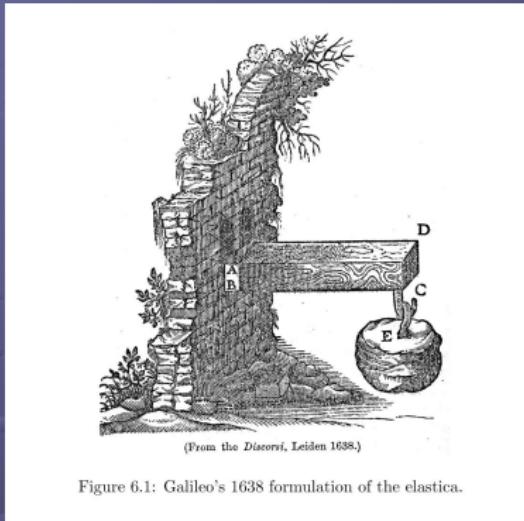
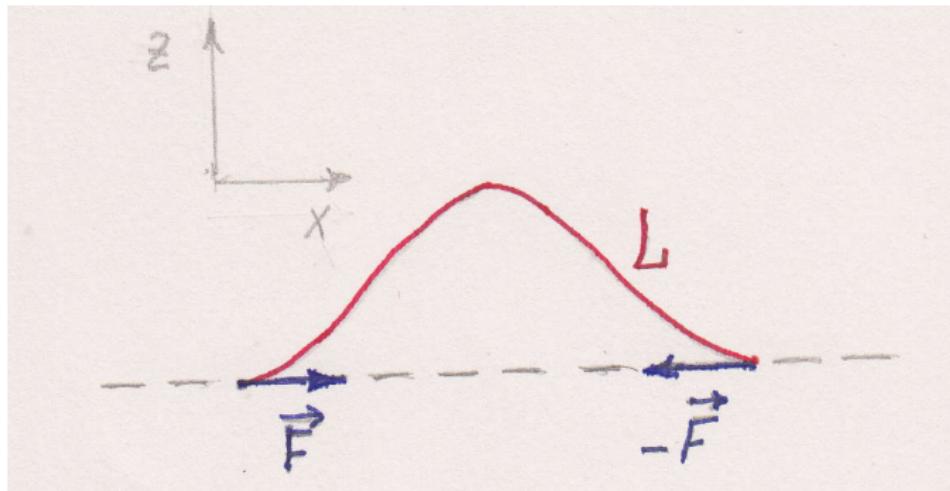


Figure 6.1: Galileo's 1638 formulation of the elastica.

Euler Elastica

Physical Prerequisites

Bending of an Elastic Rod with Fixed Length L
(elastica with tension)



Physics of the Elasticity of Rods I

Hook's Law (1660): When a thin elastic rod is bent every element of it undergoes a small strain ϵ that is in the same proportions with the stress σ in the rod

$$\sigma = E\epsilon$$

where E is the Young's modulus.

Physics of the Elasticity of Rods II

Daniel Bernoulli's Suggestion (1742): The work done in bending of an elastic rod (the energy of bending U) is proportional to the square of the curvature κ

$$U = IE \int_0^L \kappa^2(s) ds$$

where E is the Young's modulus and I is the moment of inertia of the cross-section about the neutral axis.

Mathematical Problem

Given a thin elastic rod **what will be the shape** of the rod when it is held by forces applied at its ends only?

Calculus of Variations Problem

Find the shape of the curve that minimizes the functional

$$J = \int_0^L \kappa^2(s) ds$$

subject to the constraint of a fixed length of the curve

$$L = \text{const.}$$

Fictitious Dynamical System

(Djondjorov, Hadzhilazova, Mladenov and Vassilev, 2008)

The nonlinearly coupled ordinary differential equations ($\lambda > 0$)

$$\ddot{x} - \lambda z \dot{z} = 0$$

$$\ddot{z} + \lambda z \dot{x} = 0$$

is equivalent to the equation

$$\ddot{\kappa}(s) + \frac{1}{2}\kappa^3(s) + \sigma\kappa(s) = 0$$

which is the intrinsic equation of the elastica with tension ($\mu \in \mathbb{R}$)

$$\sigma = \lambda\mu$$

Fictitious Dynamical System
(Djondjorov, Hadzhilazova, Mladenov and Vassilev, 2008)

By integrating once the fictitious dynamical system takes the form

$$\begin{aligned}\dot{x} &= \frac{\lambda z^2}{2} + \mu \\ \dot{z}^2 &= -\frac{\lambda^2 z^4}{4} - \lambda \mu z^2 - \mu^2 + 1\end{aligned}$$

where $\lambda > 0$, $\mu < 1$. In obtaining the above equations it was assumed that the particle trajectory is traced with unit speed

$$\dot{x}^2 + \dot{z}^2 = 1.$$

Alternative Parametrizations

Via the Jacobian Elliptic Functions and Elliptic Integrals

Euler Elastica for $\mu \in (-1, 1)$

(Djondjorov, Hadzhilazova, Mladenov and Vassilev, 2008)

$$x(s) = \frac{2}{\sqrt{\lambda}} E(\operatorname{am}(\sqrt{\lambda}s), k) - s, \quad z(s) = a \operatorname{cn}(\sqrt{\lambda}s, k)$$

where

$$a = \sqrt{\frac{2(1-\mu)}{\lambda}}, \quad k = \sqrt{\frac{1-\mu}{2}}$$

$E(u, k)$ incomplete elliptic integral of second order

$\operatorname{am}(u, k)$ Jacobian amplitude function

$\operatorname{cn}(u, k)$ Jacobian elliptic cosine function

Alternative Parametrizations

Via the Jacobian Elliptic Functions and Elliptic Integrals

Euler Elastica for $\mu = -1$

(Djondjorov, Hadzhilazova, Mladenov and Vassilev, 2008)

$$x(s) = \frac{2 \tanh(\sqrt{\lambda}s)}{\sqrt{\lambda}} - s, \quad z(s) = \frac{2 \operatorname{sech}(\sqrt{\lambda}s)}{\sqrt{\lambda}}$$

Alternative Parametrizations

Via the Jacobian Elliptic Functions and Elliptic Integrals

Euler Elastica for $\mu < -1$

(Djondjorov, Hadzhilazova, Mladenov and Vassilev, 2008)

$$x(s) = aE(\operatorname{am}(\sqrt{\frac{\lambda(1-\mu)}{2}}s, k), k) + \mu s$$

$$z(s) = a \operatorname{dn}(\sqrt{\frac{\lambda(1-\mu)}{2}}s, k)$$

where

$$a = \sqrt{\frac{2(1-\mu)}{\lambda}}, \quad k = \sqrt{\frac{2}{1-\mu}}$$

$E(u, k)$ incomplete elliptic integral of second order

$\operatorname{am}(u, k)$ Jacobian amplitude function

$\operatorname{dn}(u, k)$ Jacobian elliptic cosine function

Alternative Parametrizations

Via the Weierstrassian Functions

Euler Elastica via the Weierstrassian Functions (a new characterization)

$$x(s) = \frac{2}{\lambda} \left[2\zeta(s) + \frac{12\wp'(s)}{12\wp(s) - 2\lambda\mu + 3\lambda} \right] + \frac{2\mu}{3}s$$
$$z(s) = \frac{2(1-\mu)}{\lambda} \cdot \frac{12\wp(s) - 2\lambda\mu - 3\lambda}{12\wp(s) - 2\lambda\mu + 3\lambda}$$

where $\wp(s)$, $\wp'(s)$ and $\zeta(s)$ are the Weierstrassian functions

$$\wp(s) \equiv \wp(s; g_2, g_3), \quad \wp'(s) \equiv \wp'(s; g_2, g_3), \quad \zeta(s) \equiv \zeta(s; g_2, g_3)$$

with the invariants

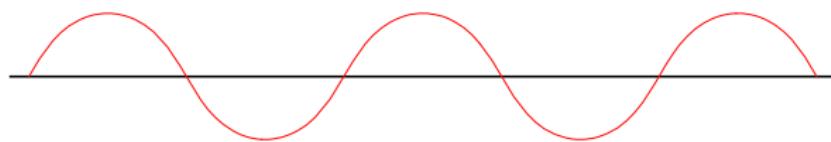
$$g_2 = \frac{\lambda^2(4\mu^2 - 3)}{12}, \quad g_3 = \frac{\lambda^3\mu(9 - 8\mu^2)}{216}$$

Alternative Parametrizations

Euler Elastica and Mathematica®

Case I

Euler Elastica for $\lambda = 4, \mu = 0.5$

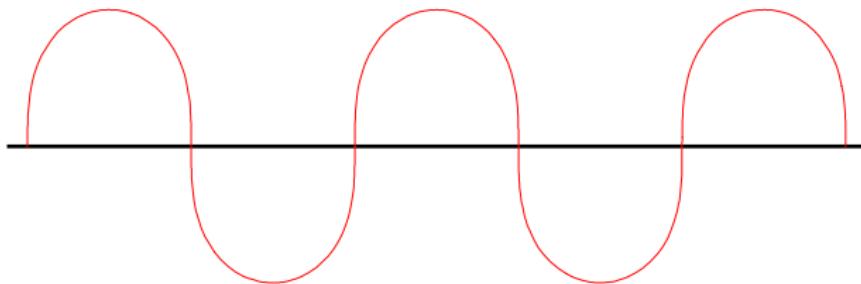


Alternative Parametrizations

Euler Elastica and Mathematica®

Case II

Euler Elastica for $\lambda = 4, \mu = 0$
(rectangular elastica)

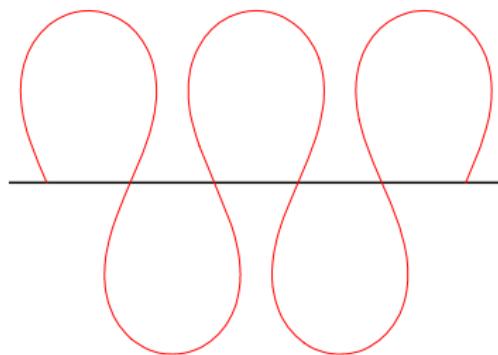


Alternative Parametrizations

Euler Elastica and Mathematica®

Case III

Euler Elastica for $\lambda = 4, \mu = -0.4$

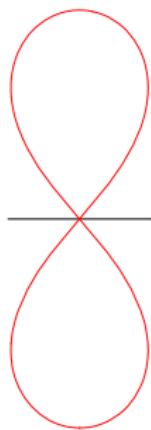


Alternative Parametrizations

Euler Elastica and Mathematica®

Case IV

Euler Elastica for $\lambda = 4$, $\mu = -0.652231$
(figure eight)

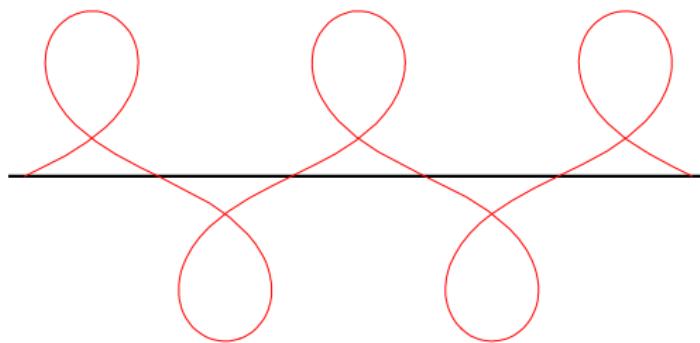


Alternative Parametrizations

Euler Elastica and Mathematica®

Case V

Euler Elastica for $\lambda = 4, \mu = -0.9$

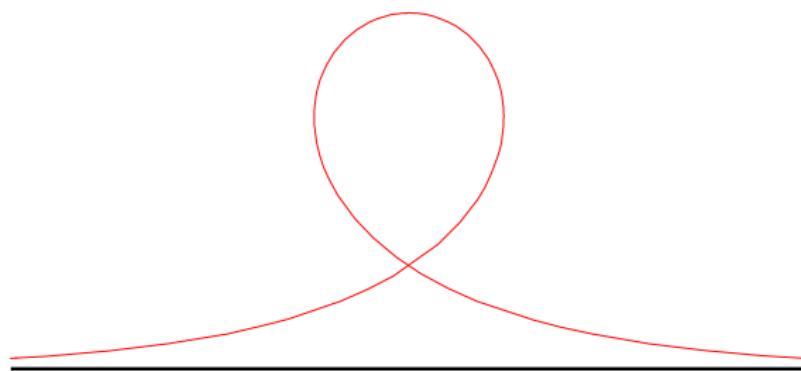


Alternative Parametrizations

Euler Elastica and Mathematica®

Case VI

Euler Elastica for $\lambda = 4, \mu = -1$

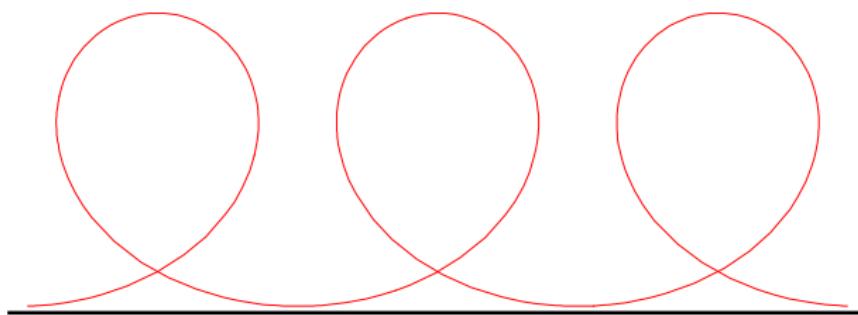


Alternative Parametrizations

Euler Elastica and Mathematica®

Case VII

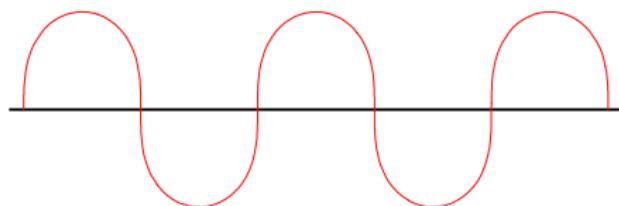
Euler Elastica for $\lambda = 4, \mu = -1.2$



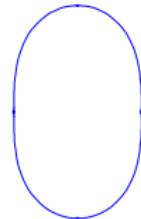
The Rectangular Elastica and the Mylar Balloon

The Mylar Balloon – A Surface of Revolution

Rectangular Elastica ($\mu = 0$)



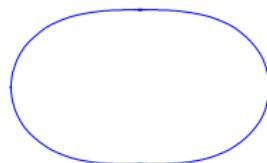
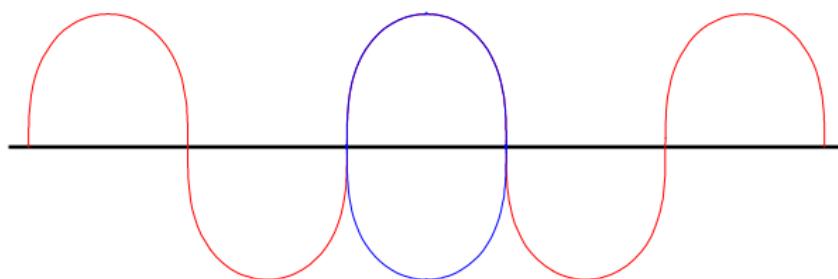
The Meridional Profile of the Mylar Balloon ($r = \sqrt{2/\lambda}$)



The Rectangular Elastica and the Mylar Balloon

The Mylar Balloon – A Surface of Revolution

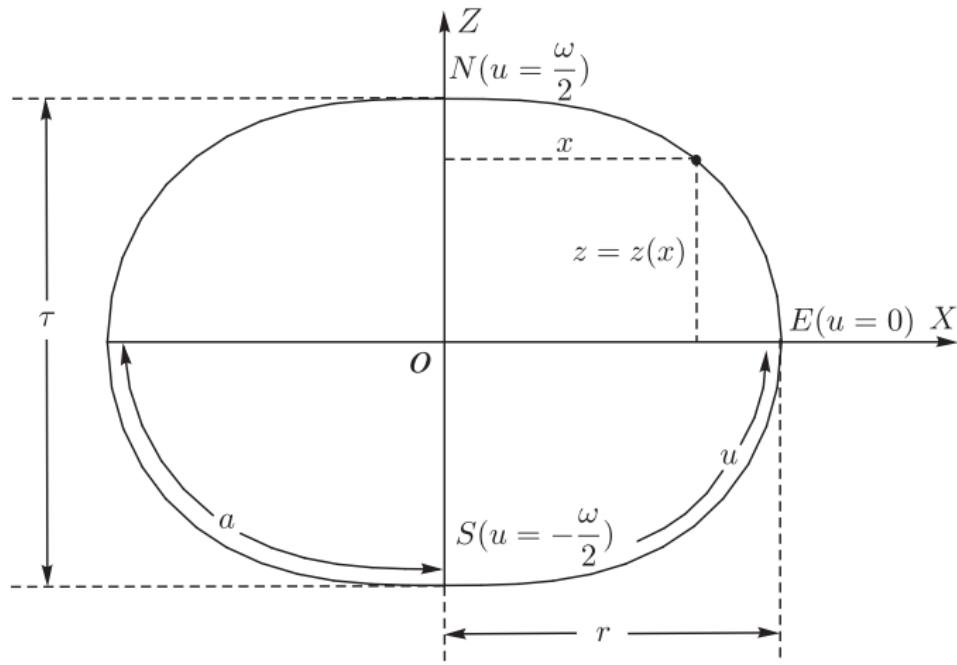
The meridional profile of the Mylar Balloon
is obtained by rotating at $\pi/2$ the Rectangular Elastica



The Rectangular Elastica and the Mylar Balloon

The Mylar Balloon – A Surface of Revolution

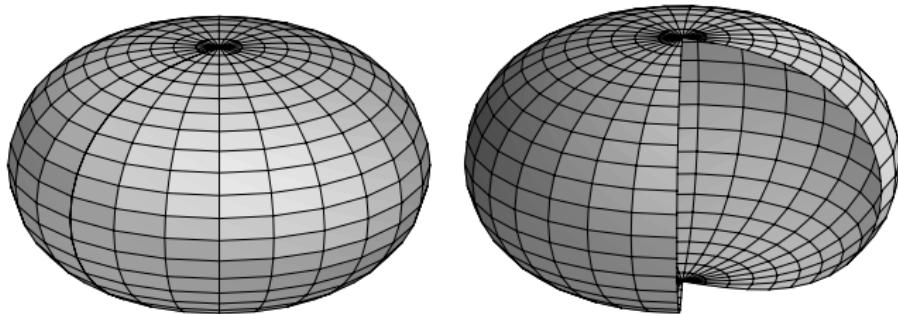
The meridional profile of the Mylar Balloon



The Rectangular Elastica and the Mylar Balloon

The Mylar Balloon – A Surface of Revolution

Two Views of the Mylar Balloon



The Rectangular Elastica and the Mylar Balloon

The Mylar Balloon – A Surface of Revolution

The Profile of the Mylar Balloon
(Pulov, Hadzhilazova and Mladenov, 2014)
 $(\mu = 0, r = \sqrt{2/\lambda})$

$$x(u) = r \frac{2\wp(u) - r^2}{2\wp(u) + r^2}$$

$$z(u) = 2\zeta(u) + \frac{2\wp'(u)}{2\wp(u) + r^2}$$

where $\wp(s)$, $\wp'(s)$ and $\zeta(s)$ are the Weierstrassian functions

$$\wp(s) \equiv \wp(s; g_2, g_3), \quad \wp'(s) \equiv \wp'(s; g_2, g_3), \quad \zeta(s) \equiv \zeta(s; g_2, g_3)$$

with the invariants

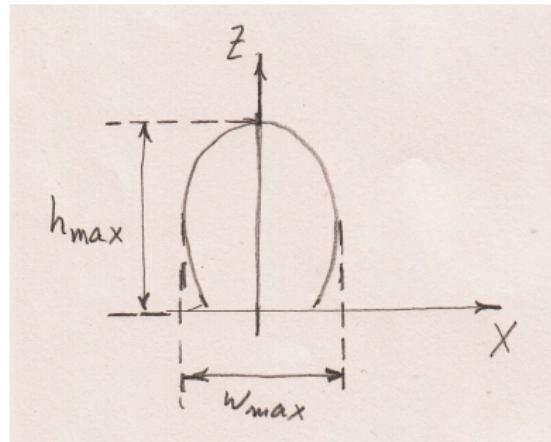
$$g_2 = -r^4, \quad g_3 = 0$$

The Rectangular Elastica and the Mylar Balloon

Aspect Ratio of the Elastica

Aspect Ratio: $\eta = h_{\max}/w_{\max}$

$$\eta(\mu) = \frac{\sqrt{1-\mu}}{\sqrt{2}(2E(\arccos\sqrt{-\mu/(1-\mu)}, \sqrt{(1-\mu)/2}) - F(\arccos\sqrt{-\mu/(1-\mu)}, \sqrt{(1-\mu)/2}))}$$



The Rectangular Elastica and the Mylar Balloon

Aspect Ratio of the Elastica

The aspect ratio of the Mylar Balloon (rectangular elastica) is

$$\eta(0) = \tilde{\omega}/\pi$$

where

$$\tilde{\omega} \approx 2.6220$$

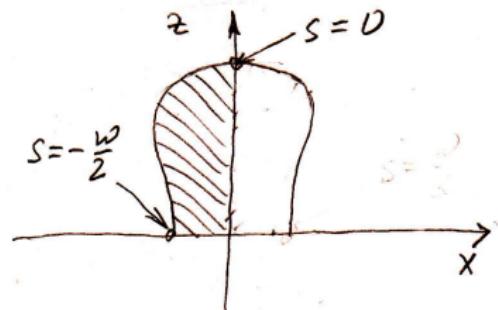
is the lemniscate constant.

The Rectangular Elastica and the Mylar Balloon

The Area Bounded by the Elastica

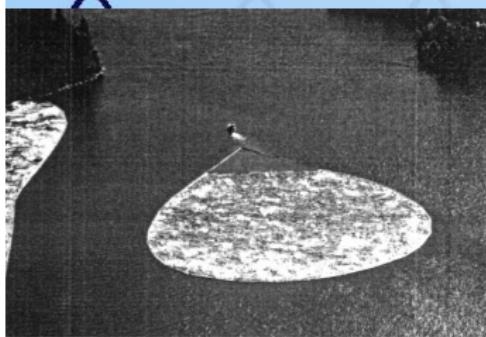
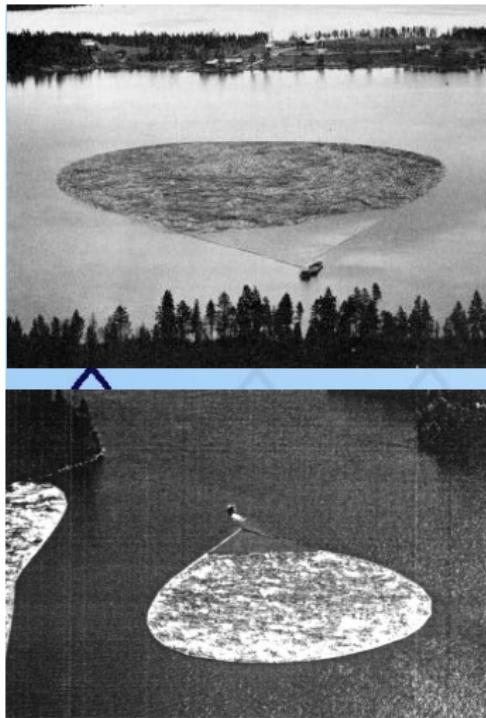
The Total Area A_{tot}
bounded by the elastica and the X axis

$$A_{\text{tot}} = 2\sqrt{1 - \mu^2}/\lambda$$



The Rectangular Elastica and the Mylar Balloon

The Area Bounded by the Elastica



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